Name and section: ____

1. The partition function, q, can generally be written as

$$q = \sum_{i} e^{-\beta \epsilon_{i}} \tag{1}$$

where ϵ_i is the energy of the state *i* and $\beta = 1/(k_B T)$.

(a) Derive the average energy in terms of the partition function. Starting from the expression

$$\langle \epsilon \rangle = \sum_{i} \epsilon_{i} p(\epsilon_{i}) = \sum_{i} \frac{\epsilon_{i} \mathrm{e}^{-\beta \epsilon_{i}}}{q}$$
(2)

where $p(\epsilon_i)$ is the probability of occupying state *i*. Derive

$$\langle \epsilon \rangle = -\frac{1}{q} \frac{\partial q}{\partial \beta} = -\frac{\partial \ln q}{\partial \beta} \tag{3}$$

using eqs. (1) and (2).

Solution: First, using the chain rule:

$$\frac{\partial \ln q}{\partial \beta} = \frac{1}{q} \frac{\partial q}{\partial \beta}$$

which shows part. Next,

$$-\frac{1}{q}\frac{\partial q}{\partial \beta} = -\frac{1}{q}\frac{\partial}{\partial \beta}\sum_{i} e^{-\beta\epsilon_{i}}$$
$$= -\frac{1}{q}\sum_{i}-\epsilon_{i}e^{-\beta\epsilon_{i}}$$
$$= \sum_{i}\frac{\epsilon_{i}e^{-\beta\epsilon_{i}}}{q}$$

(b) Now let's study the temperature dependence of the average vibrational energy of a diatomic molecule. We learned from CH351 that the vibrational motion of a diatomic molecule can be described as a one-dimensional harmonic oscillator problem. If the molecule vibrates with frequency ω , what are the quantized vibrational energy levels?

Solution: $\epsilon_i = \hbar \omega \left(i + \frac{1}{2} \right)$

(c) Let's derive the vibrational partition function for this diatomic molecule. Starting with the definition of the partition function given in eq. (1), show that

$$q = \frac{e^{-\beta\hbar\omega/2}}{1 - e^{-\beta\hbar\omega}}.$$
(4)
Solution:

$$q = \sum_{i=0}^{\infty} e^{-\beta\hbar\omega(i+\frac{1}{2})}$$

$$= e^{-\beta\hbar\omega/2} \sum_{i=0}^{\infty} e^{-\beta\hbar\omega i}$$

$$= e^{-\beta\hbar\omega/2} \frac{1}{1 - e^{-\beta\hbar\omega}}$$

$$q = \frac{e^{-\beta\hbar\omega/2}}{1 - e^{-\beta\hbar\omega}}$$

(d) Show that the average vibrational energy for the diatomic molecule is

$$\langle \epsilon \rangle = \frac{1}{2} \hbar \omega \left(\frac{1 + e^{-\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}} \right).$$
(5)

Solution:

$$\begin{split} \langle \epsilon \rangle &= -\frac{\partial}{\partial \beta} \left(-\beta \hbar \omega / 2 - \ln \left(1 - e^{-\beta \hbar \omega} \right) \right) \\ &= \frac{\hbar \omega}{2} + \frac{\hbar \omega e^{-\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}} \\ &= \frac{1}{2} \hbar \omega \left(\frac{1 - e^{-\beta \hbar \omega} + 2e^{-\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}} \right) \\ \langle \epsilon \rangle &= \frac{1}{2} \hbar \omega \left(\frac{1 + e^{-\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}} \right) \end{split}$$

(e) In the limit as $T \to 0$, to what value does $\langle \epsilon \rangle$ converge?

Solution: As $T \to 0, \beta \to \infty$:

$$\lim_{\beta \to \infty} \left< \epsilon \right> = \frac{1}{2} \hbar \omega \left(\frac{1}{1} \right)$$

as is expected because of the zero point energy for quantum harmonic oscillators.

(f) What about the limit of $T \to \infty$? Taylor series expand the exponential terms to first order in $\beta \hbar \omega$ to estimate the *classical* limit of $\langle \epsilon \rangle$.

Solution: As $T \to \infty$, $\beta \to 0$: Taking the Maclaurin Series of $e^{-\beta\hbar\omega}$, $e^{-\beta\hbar\omega} \approx 1 - \beta\hbar\omega$ $\langle \epsilon \rangle \approx \frac{1}{2}\hbar\omega \left(\frac{2+\beta\hbar\omega}{\beta\hbar\omega}\right)$ $= \frac{1}{\beta} + \frac{1}{2}\hbar\omega$ $\lim_{T\to\infty} \langle \epsilon \rangle \approx k_B T$

(g) Graphically show how $\langle \epsilon \rangle$ varies with T.

